

Random considerations on KPZ

The "P" of KPZ

Antefacts

I was fascinated by DLA (Diffusion Limited aggregation). Apparently no upper critical dimensions.

With Zhang Yicheng we tried to study a simpler problem. Aggregation: i.e. the Eden model (*Eden model in many dimensions*). We computed an $1/d$ expansion, nothing interesting was found.

We wrote a paper *Field theories and growth models: Eden model and DLA are related to Reggeon field* but we were not able to make progress.

The Eden model has surface fluctuations. We could study them, but Zhang left in 1984 for Brookheaven.



$$\dot{\varphi} = V - \Delta \varphi + \eta(x, \varphi)$$

$$\dot{\varphi} = +\Delta \varphi + (\partial_{\mu} \varphi)^2 + \eta(x, \varphi)$$

1

$$\phi = V + i$$

$$\vec{p} = V - \Delta\phi + \eta(x, \phi)$$

Koplink

$$\varphi = Vc + \tilde{\varphi}$$

$$g \sim \frac{1}{V}$$

$$\tilde{\varphi} = \varphi + \Delta\phi + \eta(x, t + g\phi)$$

g_c

$$+ \eta(t, c) + g\phi \partial_t \eta$$

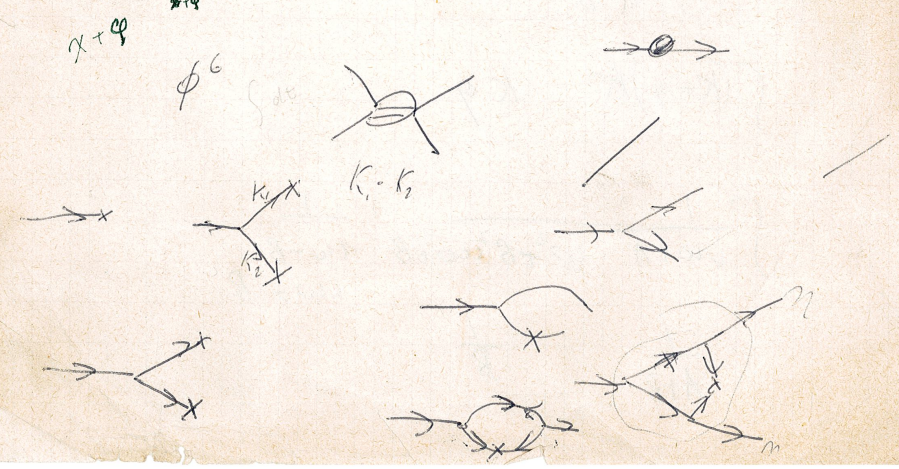
$$\tilde{\varphi} = +\Delta\phi + (\partial_{\mu}\phi)^2 + \eta(x, c)$$

$$[\varphi] = \frac{D}{2} - 1$$

$$[\varphi] = 0$$

$$L(\partial_{\mu}\phi) =$$

$$[D=2]$$



Interface moving through a random background

Joel Koplik and Herbert Levine

Schrödinger equation of zero interacting particles

KPZ \rightarrow polymer in Random Potential \rightarrow Schrödinger equation of n interacting particles in the limit $n \rightarrow 0$.

$$\mathbf{H} = -\frac{1}{2} \sum_{a=1,n} \Delta_a - \sum_{a=1,n,b=1,n} V(x_a - x_b),$$

$$\overline{\rho(x_1)\rho(x_2)} \equiv \int dP[\rho] \rho(x_1)\rho(x_2) \propto \int dx_3 \dots dx_n \psi(x_1, \dots, x_n).$$

Wonderful paper of Kardar in $d = 1$: the solution with a δ potential is known for any n :

$$\Phi(x_1 \cdots x_n) = \prod_{i,k} \psi(x_i - x_k)$$

A variational approach to directed polymers (Mézard, GP)

$$\psi(x_1, \dots, x_n) \propto \int dR[\mu] \mu(x_1), \dots, \mu(x_n).$$

$$\langle \psi | \mathbf{H} | \psi \rangle = \int dR[\mu] dR[\sigma] \mathfrak{S}_n[\mu, \sigma]$$

$$\begin{aligned} \mathfrak{S}_n[\mu, \sigma] = & n/2 \int dx \sum_{\nu=1, N} d\mu/dx_\nu d\sigma/dx_\nu I[\mu\sigma]^{n-1} \\ & - n(n-1) \int dx dt \mu(x)\mu(y)\sigma(x)\sigma(y)V(x-y)I[\mu\sigma]^{n-2} \end{aligned}$$

$$\langle \psi | \psi \rangle = \int dR[\mu] dR[\sigma] I[\mu\sigma]^n.$$

$$\int dR[\sigma] \mathfrak{S}_n[\mu, \sigma] = E_n \int dR[\sigma] I[\mu\sigma]^n$$

Schrödinger eq.

Half baked ideas.

The paper is full of half baked ideas.

Greens functions identities, replica symmetry breaking....

Hard problem:

$$\psi(x_1, \dots, x_n) \propto \int dR[\mu] \mu(x_1), \dots, \mu(x_n).$$

Do we have an $R_n[\mu]$ that is independent from n . Is it unique?

RSB:

On the 1/D expansion for directed polymers with Slanina.

$$\exp \left(-\frac{1}{4} \sum_{a,b} Q_{ab}^{-1} \omega_a \cdot \omega_b \right)$$

A desperate attempt to do a too difficult computation using educated guesses.

I think that now we could do the computation without educated guesses.

Numerical Simulations: RSOS

Marinari, Pagnani, GP.

$$h(x+1) - h(x) = \{-1, 0, 1\}$$

The problem is formulated on link variables (like gauge fields). The field should be pure gauge at time 0.

We random choose the point x and we increase all the links starting from x if this is possible.

Gauge fields are represented as two bits. One reconstruct the surface when needed for measurements.

We look to the limit $t \rightarrow \infty$: the width interface w_2 scales as L^χ .

RSOS

If one looks to the h 's, RSOS has rules similar to self-organized criticality (sandpiles).

Discrete space, dynamical rules. Random cellular automata with few states per site.

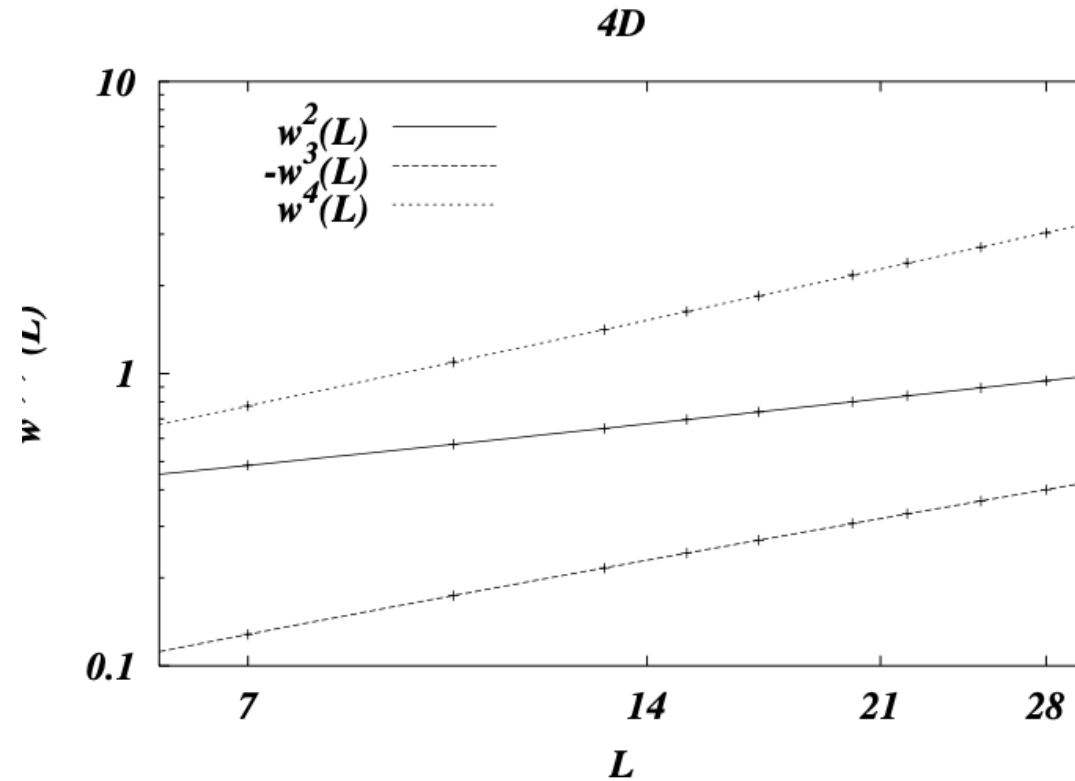
The h 's have long range anticorrelations that have a power decay.

How large is the universality class?

Are there other models with these features that are not trivial in high dimensions?

$D = 4$ $L = 7 \cdots 28$: L^4 points

We study the cumulant of $h(x, t) - h(t)$: w_2, w_3, w_4



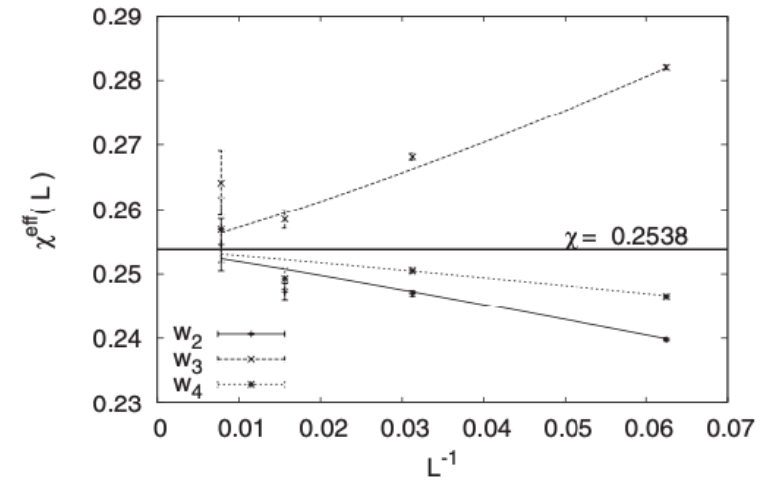
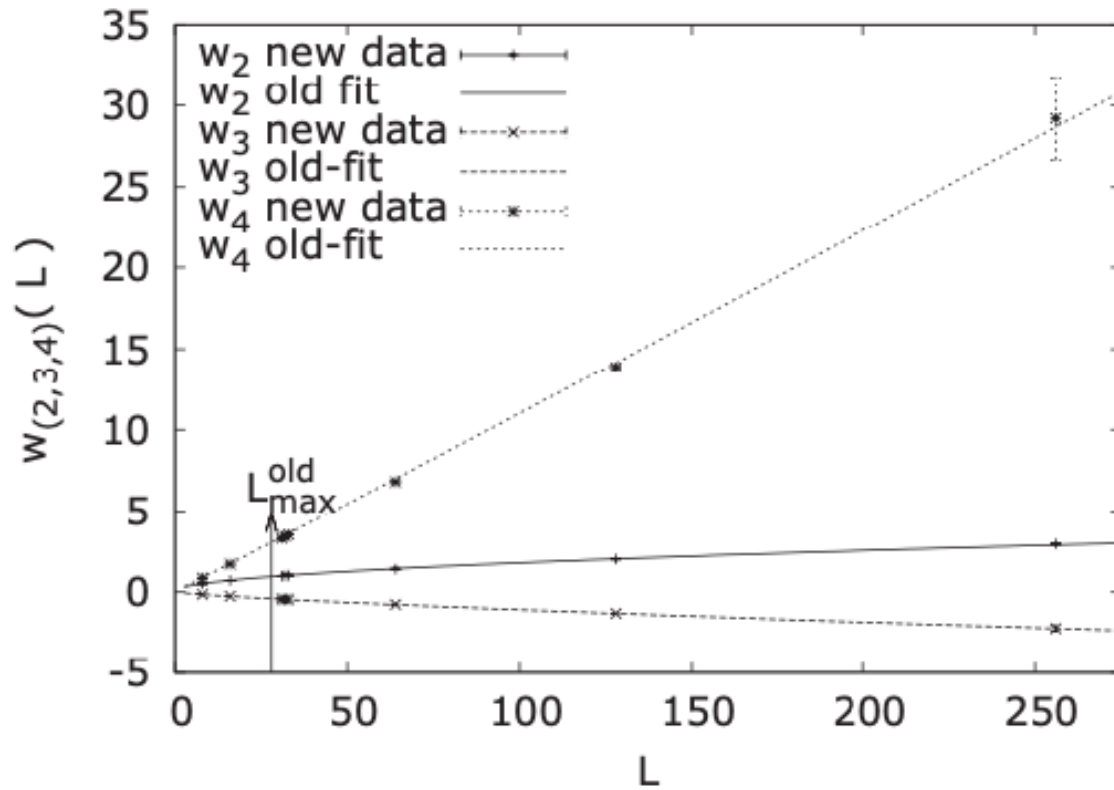
$$w_2 \simeq A_2 L^{2\chi} (1 + B_2 L^{-\omega})$$

$$w_3 \simeq -A_3 L^{3\chi} (1 + B_3 L^{-\omega})$$

$$w_4 \simeq A_4 L^{4\chi} (1 + B_4 L^{-\omega}).$$

$$\chi_{D=4} = 0.255 \pm 0.003 \quad \omega_{D=4} = 0.98 \pm 0.09.$$

Vingt ans après: $D = 4$ $L = 7 \cdots 128(256)$



χ and ω

New	0.2537(8)	1.11(9)
Old	0.255(3)	0.98(9)

Vingt ans après: $D = 2$ same L range with much higher statistics

$$w_2 = A_2 L^{2\chi} (1 + B_2 L^{-\omega}),$$

$$w_3 = S A_2^{3/2} L^{3\chi} (1 + B_3 L^{-\omega}),$$

$$w_4 = K A_2^2 L^{4\chi} (1 + B_4 L^{-\omega}).$$

$$w_2 = A_2 L^{2\chi} (1 + B_2 L^{-\omega} + C_2 L^{-2\omega}),$$

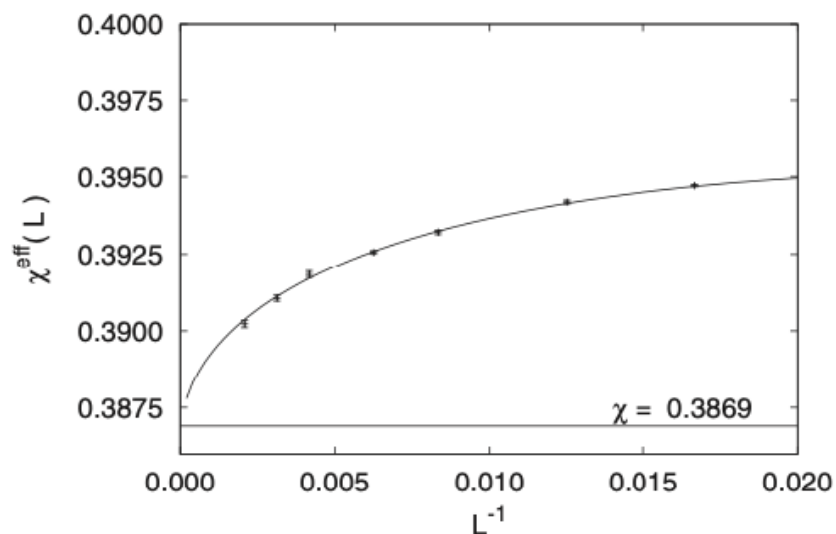
$$w_3 = S A_2^{3/2} L^{3\chi} (1 + B_3 L^{-\omega} + C_3 L^{-2\omega}),$$

$$w_4 = K A_2^2 L^{4\chi} (1 + B_4 L^{-\omega} + C_4 L^{-2\omega}).$$

	χ	ω	A_2	B_2	C_2
FIT I	0.3893(6)	0.8(2)	0.118(1)	-0.4(2)	NA
FIT II	0.3869(4)	0.57(5)	0.1226(1)	-0.37(2)	0.6(2)

$$\chi_{D=2} = 0.393 \pm 0.003$$

$$\omega_{D=2} = 1.1 \pm 0.3.$$



Width distributions

Marinari, Pagnani, GP, Racz

$$w_2 = \frac{1}{A_L} \sum_{\mathbf{r}} [h(\mathbf{r}, t) - \bar{h}]^2,$$

$$P_L(w_2) \approx \frac{1}{\langle w_2 \rangle_L} \Phi_d \left(\frac{w_2}{\langle w_2 \rangle_L} \right)$$

