



## Back to the rainbow

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<u>Thm</u>: The  $ASEP_{t,q}$  sheet converges to the **Airy sheet** S(x;y).

## Key ideas behind the theorem

- ✓ Work with colored stochastic six vertex (S6V) model
- Lecture 3 focuses on the uncolored model
  - ✓ Yang-Baxter embeds S6V height function in q-Boson model
  - ✓ Gibbs property, one-point GUE Tracy-Widom asymptotics and strong characterization yields Airy line ensemble limit
- Lecture 4 returns to the colored model
  - Yang-Baxter story extends to colored S6V, q-Boson models
  - Intercolor Gibbs property yields approximate variational representation of colored line ensemble via uncolored one
  - Yields  $ASEP \rightarrow Airy$  sheet limit via Airy line ensemble limit

## Colored stochastic six vertex (S6V) model

[Kulish-Reshetikhin-Sklyanin '81] [Bazhanov '85], [Jimbo '86], [Kuniba-Mangazeev-Maruyama-Okado '16], [Borodin-Wheeler '18]





'Quantum' parameter q 'Spectral' parameter z ASEP is recovered around the diagonal as  $z \rightarrow 1$ 



## Yang-Baxter for colored S6V and q-Boson weights





S6V weights

q-Boson weights



Yang-Baxter equation

#### Aside: Yang-Baxter RRR relations for colored S6V





#### **Aside: Fusion**

[Kulish-Reshetikhin-Sklyanin '81]

$$R_{y/x}\Big((i_1, \dots, i_M), j; (k_1, \dots, k_M), \ell\Big) := \sum_{c_1=0}^n \cdots \sum_{c_{M-1}=0}^n R_{q^{M-1}y/x}(i_1, j; k_1, c_1)R_{q^{M-2}y/x}(i_2, c_1; k_2, c_2) \dots R_{y/x}(i_M, c_{M-1}; k_M, \ell)$$

$$= \int_{i_1}^{k_1} \int_{i_2}^{k_2} \cdots \int_{i_M}^{k_M} \ell$$

$$\mathcal{L}_{y/x}^{(M)}(\mathbf{I}, j; \mathbf{K}, \ell) := \frac{1}{Z_q(M; \mathbf{I})} \sum_{\substack{C(i_1, \dots, i_M) = \mathbf{I} \\ C(k_1, \dots, k_M) = \mathbf{K}}} q^{\operatorname{inv}(i_1, \dots, i_M)} R_{y/x}\Big((i_1, \dots, i_M), j; (k_1, \dots, k_M), \ell\Big)$$

#### Aside: Fusion



## Yang-Baxter for colored S6V and q-Boson weights





S6V weights

q-Boson weights



Yang-Baxter equation

#### Relating the colored q-Boson and S6V models







#### Colored q-Boson model as a line ensemble





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#### Intercolor q-Boson Gibbs property



satisfies a variational formula at q=0:  $L_k^{(2)} = \mathsf{PT}\left(L_k^{(1)}, L_{k+1}^{(2)}\right)$ 

$$\mathsf{PT}(f,g)(y) := f(y) + \max_{0 \le y' \le y} \left( g(y') - f(y') \right)$$

and an approximate variational formula at q>0

$$\mathbb{P}\left(\max_{y\in[\![0,N+M]\!]} \left(L_k^{(2)}(y) - \mathsf{PT}(L_k^{(1)}, L_{k+1}^{(2)})(y)\right) \ge m\right) \le q^{cm^2}$$

# Origin of intercolor q-Boson Gibbs property when

1

 $i \longrightarrow j$ 

 $u(1-q^{A_j})q^{A_{[j+1,N]}}$ 

When q=0, the highest arrow incoming to each vertex must exit right.

1

 $u(1-q^{A_i})q^{A_{[i+1,N]}}$ 

When q>0 there is a probabilistic penalty.



 $uq^{A_{[i+1,N]}}$ 

 $j \longrightarrow i$ 

0

#### Iterating intercolor Gibbs property yields (q=0)



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#### Putting it all together



### Take home messages

- A full KPZ scaling limit may be overkill from a physics perspective, yet it implies the scaling limits for all sorts of observables of physical interest as corollaries.
- There is enhanced mathematical structure present in analysis of the full object compared to observables.
- We embed the full object into a larger object using the Yang-Baxter equation and extract the scaling limit of this.